

# PCMI GMT Lecture #3

continuing with 4.5.9

(15)  $E$  is the jump set, so for WCE

$\|\partial T\|(\omega) = \text{integral of height of jump along } \omega$

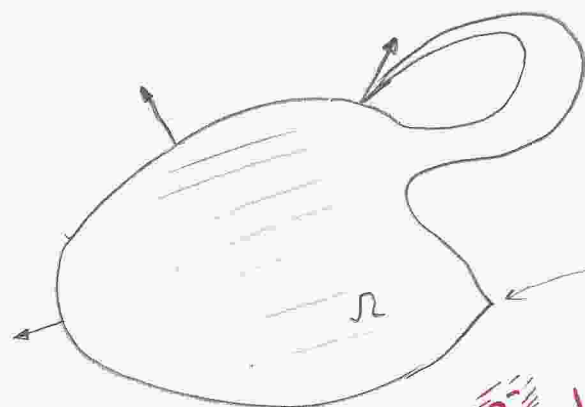
$$= \int_{\omega} (m - \lambda) \, d\mathcal{H}^{n-1}$$

(16) & (17)  $E$  is an  $n-1$  rectifiable current in  $\mathbb{R}^n$

in the case that  $f = \chi_{\Omega}$ , this means the set  $\{x \in \mathbb{R}^n : \lambda(x) < m(x)\} = \partial^* \Omega$

= a "nice" set with exterior normals at  $\mathcal{H}^{n-1}$  almost every point in  $E$

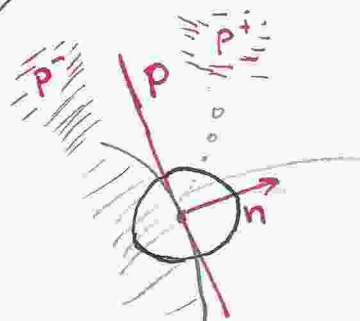
↑ reduced boundary of  $\Omega$



no exterior normal!

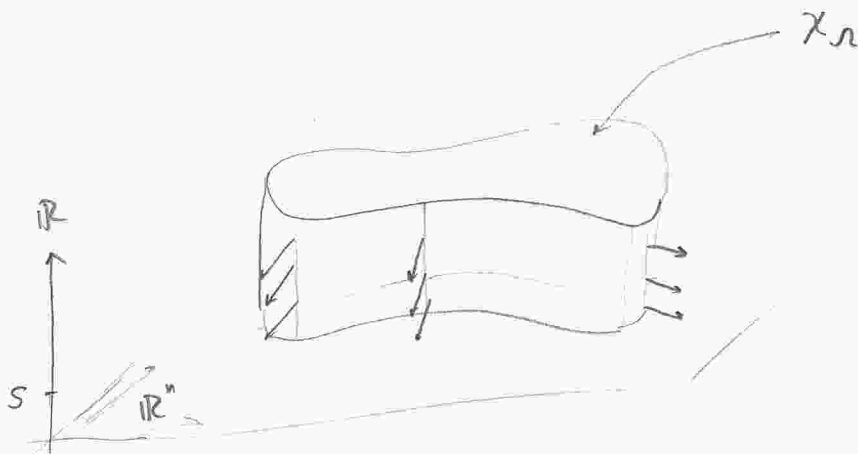
Def: Exterior normal

$P$  the hyperplane with normal  $n$  divides  $\Omega$  locally, measure theoretically into inside & outside



$$\lim_{r \rightarrow 0} \frac{\mathcal{L}^n(B_r(x) \cap P^+ \cap \Omega)}{\alpha(n) r^n} = 0$$

$$\lim_{r \rightarrow 0} \frac{\mathcal{L}^n(B_r(x) \cap P^- \cap \Omega)}{\alpha(n) r^n} = 1$$



subset of  $\mathbb{R}^n$

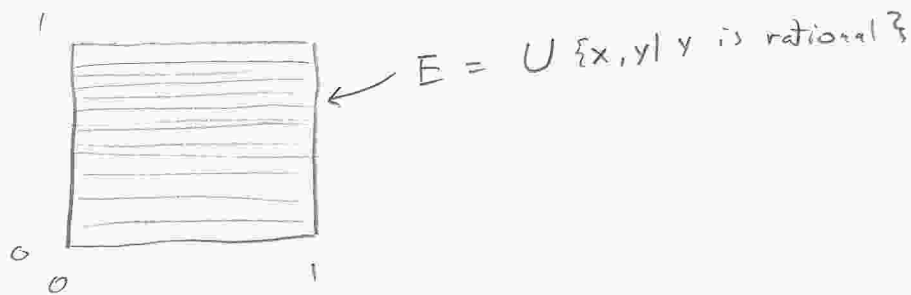
$$n[G, (\bar{b}, s)] = (n[\{x \mid f(x) \geq s\}, \bar{b}], 0)$$

vector in  $\mathbb{R}^n$

point in  $\mathbb{R}^{n+1}$

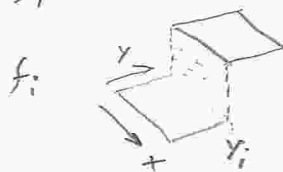
vector in  $\mathbb{R}^{n+1}$

Note:  $E$  is countably  $(\mathbb{H}^{n-1}, n-1)$ -rectifiable not  $(\mathbb{H}^{n-1}, n-1)$ -rectifiable. We can create an example  $\exists$  for any compact  $K$   $\mathbb{H}^{n-1}(K \cap E) = \infty$ .



To get an  $f$  with this jump set, enumerate the rational  $y$ ,  $\{y_1, y_2, y_3, \dots\}$  and then let  $f_i$  be defined by

$$f_i : [0, 1]^2 \rightarrow [0, 1] : (x, y) \rightarrow \frac{1}{2^i} \begin{cases} 0 & y < y_i \\ 1 & y_i \leq y \end{cases}$$



Then  $f(x, y) = \sum_i f_i(x, y)$ .

For this  $f \int_{[0,1]^2} |f| dx < 1$

and

$$\int_{[0,1]^2} |\nabla f| dx = 1$$

Remark: we can always use a real multiplicity to turn a countably  $(\mathbb{H}^k, \kappa)$ -rectifiable set into a real rectifiable current. In fact,

Defining  $T \equiv E \llcorner f$ , for the  $f$  just defined,  $\partial T = [E]$ ,  $\vec{E} = \vec{e}_1$ ,  $\rho(x, y) = \sum \underbrace{\frac{1}{2}}_{\text{multiplicity}} \cdot \mathbb{1}_{\{i \neq y_i \leq y\}}$

(19)

The condition

$$\mathcal{L}^n [U(b, \rho) \cap \{x: f(x) > t\}] \leq \alpha(n) \rho^{n/2}$$

$$\mathcal{L}^n [U(b, \rho) \cap \{x: f(x) < t\}] \leq \alpha(n) \rho^{n/2}$$

open ball, radius  $\rho$   
center  $b$

is a solution of a median value of  $f$  on  $U(b, \rho)$ .

Notice:

$$\{x: f(x) \geq t\}$$

$$\{x: f(x) \leq t\}$$

will not work in the definition.

Example ...



... both expressions yield sets with measure  $> 1/2$

of course there is an easy fix

$$\mathcal{L}^n(\dots \geq \epsilon\} \geq \alpha(n) \rho^{n/2}$$

$$\mathcal{L}^n(\dots \leq \epsilon\} \geq \alpha(n) \rho^{n/2}.$$

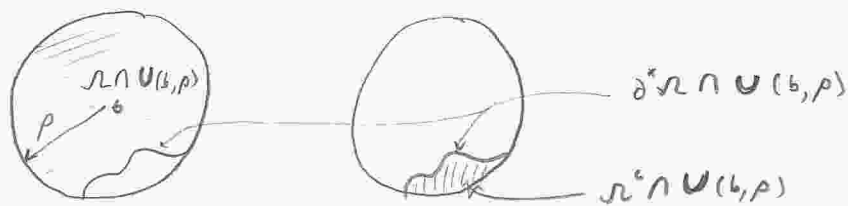
$$*) \left( \rho^{-n} \int_{U(b, \rho)} |f - \epsilon|^p d\mathcal{L}^n \right)^{1/p} \leq \sigma \rho^{1-n} \|\partial T\| (U(b, \rho)) \quad \beta = \frac{n}{n-1}$$

is Poincaré's inequality for top dimensional normal currents - i.e. Functions of Bounded variation (Everything is local here).

$\sigma$  is the relative isoperimetric ratio.

$$**) \left[ \min \{ \mathcal{L}^n(U(b, \rho) \cap \Omega), \mathcal{L}^n(U(b, \rho) \cap \Omega^c) \} \right]^{1/p} \leq \sigma \mathcal{H}^{n-1}(\partial^* \Omega \cap U(b, \rho))$$

Relative isoperimetric constant.



without the "min" we cannot get a relative isop. ratio.

\*) has \*\*) as a special case: let  $f$  be the characteristic function of  $\Omega$ ,  $\chi_\Omega$ .

Notice that it then automatically makes the left hand side of \*) into

$$\left( \rho^{-n} \min \{ \mathcal{L}^n(U(b, \rho) \cap \Omega), \mathcal{L}^n(U(b, \rho) \cap \Omega^c) \} \right)^{1/p}$$

and, after factoring out the  $\rho^{1-n}$  from both sides we get \*\*.

From (15) and the fact that  $\|\partial T\|$  is locally finite, we have that  $\|\partial T\|$  is Radon. From this we know that

$$\lim_{r \rightarrow 0} \frac{\|\partial T\|(B(b, \rho))}{\alpha(n-1) \rho^{n-1}} = 0$$

For  $\mathcal{H}^{n-1}$  a.e.  $b \in \mathbb{R}^n \setminus E$ . For such  $b$ , then, we have,  $\theta^{*n-1}(\|\partial T\|, b) = \theta_*^{n-1}(\|\partial T\|, b) = \theta^{n-1}(\|\partial T\|, b) = 0$

and

$$\lim_{\rho \rightarrow 0} \int_{U(b, \rho)} |f(x) - t|^p d\mathcal{L}^n_x = 0$$

But  $t$  depends both on  $b$  and on  $\rho$ .

(20) and (21) are two inequalities in this direction.

in 20  $\lambda(b) = \mu(b) = F(b)$  and  $F$  is approx. cont. there. But we are not guaranteed a 0 on the R.H.S. This seems a bit strange at first but ... example

① if  $f$  is essentially bounded then

$$\rho^{-n} \int_{U(b, \rho)} |f(x) - F(b)|^p d\mathcal{L}^n_x \Rightarrow 0$$

as  $\rho \rightarrow 0$  (By the definition of  $\lambda, \mu$  and  $F$ )

② Note that if  $b=0$ ,  $n=2$  and  $f$  is radially symmetric with

$$f(r) = \begin{cases} \frac{1}{\sqrt{r}} & 0 \leq r \leq L \\ 0 & r > L \end{cases}$$

$$\int |f| < \infty \quad \text{and} \quad \int |\nabla f| < \infty$$

but  $F(0) = \infty \Rightarrow$  LHS of the inequality in (20) is  $\infty$ . computes:

$$\int_0^{\rho} \int_0^{2\pi} \frac{1/2}{r^{3/2}} d\theta dr = \pi \int_0^{\rho} r^{-1/2} dr = 2\pi r^{1/2} \Big|_0^{\rho} = 2\pi \rho^{1/2}$$

$$\Rightarrow \text{RHS} = \sigma \alpha(1) \limsup_{\rho \rightarrow 0} \frac{2\pi \rho^{1/2}}{\alpha(1) \rho^1} = \sigma \cdot \infty$$

Putting every thing together

- (1)  $\lambda = \mu = F$  on  $\mathbb{R}^n \setminus E$
- (2)  $\theta^{n-1}(\|\partial T\|, b) = 0$   $\mathcal{H}^{n-1}$  a.e. in  $\mathbb{R}^n \setminus E$
- (3) (20) (notice that (20) was an everywhere not an a.e. statement)

implies

(21)

(22) measure theoretic normal again.

(29) and (30) when comparing these similar sets of conditions, 29-III and 30-V is one place to see what is different between (29) & (30).

(29) characterizes the jump set part of  $\|\partial T\|$  and (30) characterized the center like part of  $\|\partial T\|$ . These two pieces contain all of  $\|\partial T\|$  that is singular w.r.t  $\mathbb{L}^n$ .